



## Communication

# A laser synthesis of vanadium oxide bonded graphene for high-rate supercapacitors

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## ABSTRACT

Graphene is a type of promising electrode material for high-energy and high-power density supercapacitors, but its electrochemical performance is greatly limited by the restacking problem. In this work, we reported a facile approach to synthesis graphene with chemically bonded vanadium oxide (VO<sub>x</sub>) nanoparticles and demonstrated that chemically-bonded VO<sub>x</sub> nanoparticles can effectively prevent the graphene sheets from restacking and hence improve the electrochemical performance. The capacitance of VO<sub>x</sub>-bonded graphene increases to 272 F/g compared to 183 F/g of pristine graphene in 1 M H<sub>3</sub>PO<sub>4</sub> aqueous electrolyte at 2 A/g. The VO<sub>x</sub>-bonded graphene also showed improved rate capability in both H<sub>3</sub>PO<sub>4</sub> and ionic liquid electrolytes. The capacitance retention increases to 54.5% from 28.5% at 100 A/g (compare to 2 A/g) in H<sub>3</sub>PO<sub>4</sub> and increases to 65.1% from 46.3% at 2 A/g (compare to 0.2 A/g) in neat ionic liquid. A high energy density of 84.4 Wh/kg is obtained within the voltage window of 4 V in ionic liquid. Even at a high-power density of 1000 W/kg, the VO<sub>x</sub>-bonded graphene shows a high energy density of 47.3 Wh/kg.

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Supercapacitors have been intensively studied as electrochemical energy storage devices due to their ultrafast charge/discharge capability and ultra-long cycle life [1–6]. As a type of most widely commercialized supercapacitors, electrical double layer capacitors (EDLCs), which store/deliver electrical energy via physical adsorption/desorption of electrolyte ions on the electrode/electrolyte interfaces, show outstanding fast energy storage/delivery performance [7–10]. However, due to the mechanism of surface energy storage, EDLCs exhibit ultra-low energy density. Increasing the operate voltage window is an effective way to increase the energy density of EDLCs by using organic/ionic liquid electrolytes, according to the relationship of  $E = 1/2 CV^2$ , where  $E$  is the energy density,  $C$  is the capacitance and  $V$  is the working voltage window. However, the power density will decrease dramatically when organic/ionic liquid electrolytes are used due to the low ion conductivity and large ion size of the electrolyte ions [11,12]. As a result, the application of EDLCs has been greatly limited since in most of the energy storage applications, both high-energy and high-power density are required. This problem can be solved by either increase the capacitance in aqueous electrolytes or improve the rate performance in organic/ionic liquid electrolytes.

Graphene is a promising candidate for high-energy and high-power density supercapacitors because of its high surface area (high capacitance, contributing to energy density) and good conductivity (high rate performance, contributing to power density) [13–17]. The intrinsic capacitance of single-layer graphene was reported to be  $\sim 21 \mu\text{F}/\text{cm}^2$  and the theoretical maximum capacitance for EDLC based on graphene materials could, in principle, reach  $\sim 550 \text{ F/g}$  if the entire surface area can be used [18]. However, the current practical graphene-based capacitances are generally lower than 100 F/g due to the restacking problem [19–24]. The 2D graphene oxide (GO)/graphene sheets can easily stack together and overlap, which results in the loss of effective surface area and hinders the vertical ion transportation. Both the capacitance and the rate capability of graphene electrode are limited by the restacking problem, which results in low energy and low power density. There are various ways in solving the restacking problem of graphene film [25,26]. Luo et al. showed that the capacitance of restacked graphene was as low as 36 F/g when charged/discharged at a current density of 2 A/g in aqueous potassium hydroxide [27]. They successfully synthesized crumpled graphene balls to alleviate the restacking problem and achieved a much higher capacitance of 118 F/g at the same current density. El-Kady developed a facile laser reduction method and achieved a much higher capacitance value of 202 F/g in H<sub>3</sub>PO<sub>4</sub> aqueous electrolyte at 1 A/g

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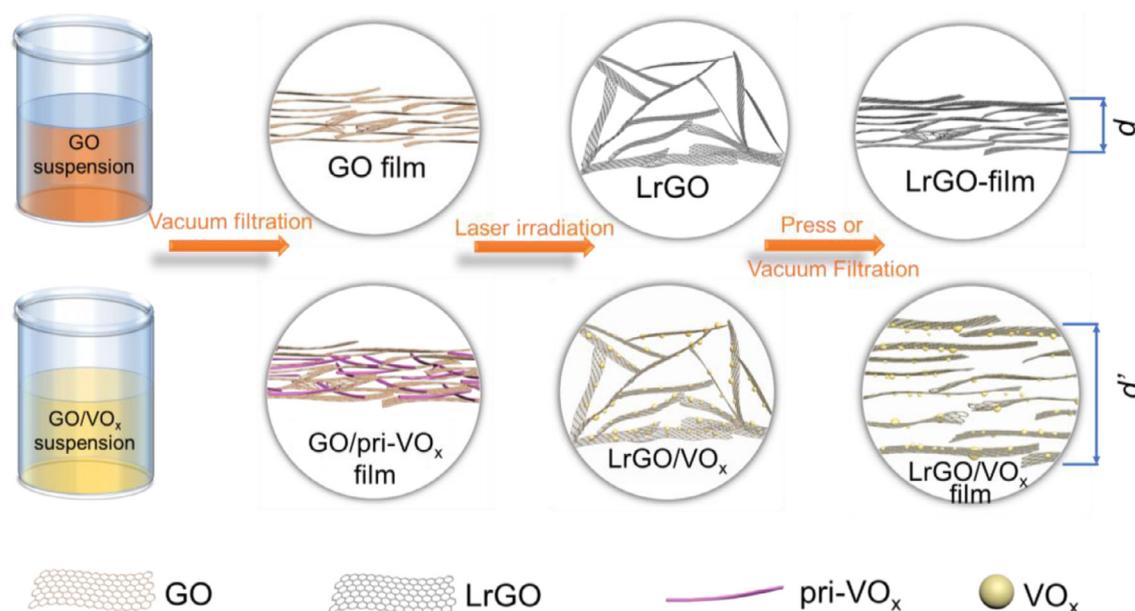


Fig. 1. Schematic illustration of laser reduced graphene and  $\text{VO}_x$  bonded graphene processes.

[28]. Kaner's work demonstrated that the GO film can be expanded by laser irradiation, which converts the restacked GO sheets immediately into well-exfoliated few-layered grapheme film, resulting in high capacitance. Luo's and El-Kady's works represent two typical strategies in solving the restacking problem of graphene sheets, which could be concluded as unstacking treatments pre- and post-GO/graphene film forming (pretreatment and posttreatment) method, respectively. Various works, afterwards, have further proved and demonstrated the effectiveness of pretreatment or posttreatment methods [26,29–32]. However, the curved nature of graphene by Luo's method and the expansion of graphene by El-Kady's laser method will largely decrease the volumetric energy density. Direct press of the curved/expanded graphene can directly increase the volumetric energy density but will cause the restacking problem again [33].

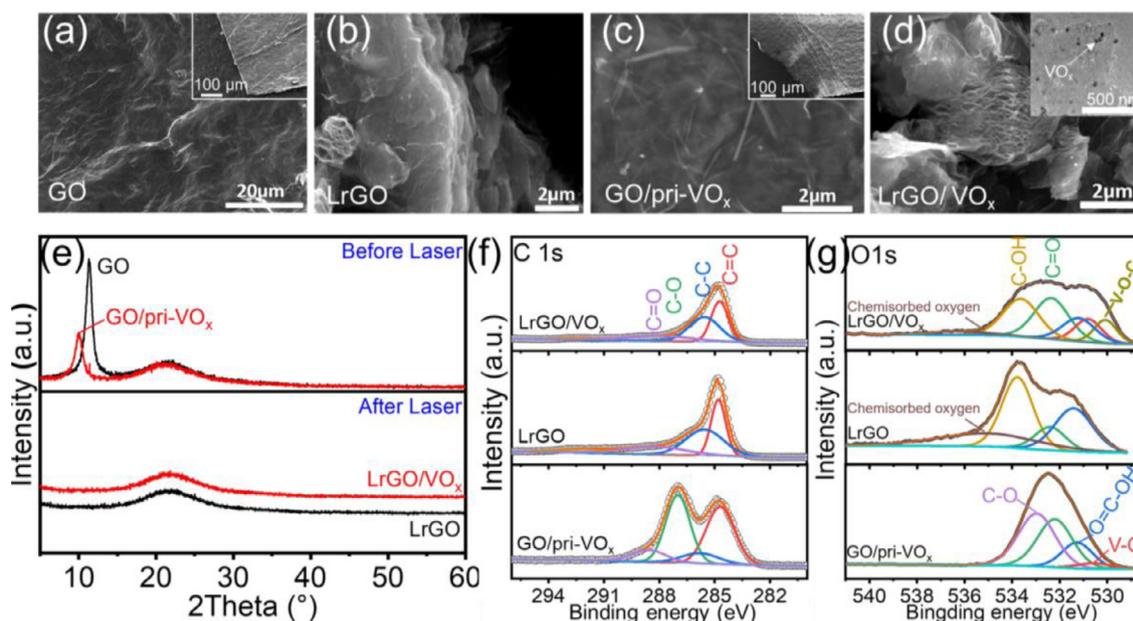
Herein, we integrated both pretreatment and posttreatment methods for the first time to further solve the restacking problem. In this work, we showed that pristine vanadium oxide ( $\text{pri-VO}_x$ ) nanofibers inserted between GO sheets can protect GO film from restacking together. Afterward,  $\text{pri-VO}_x$  nanofibers were laser converted into nanoparticles which are chemically bonded to the graphene surface. Synchronously, the separated GO sheets were further laser exfoliated and reduced. Accordingly,  $\text{VO}_x$  nanoparticles effectively prevent the graphene sheets from restacking when press or vacuum filtrate the loose reduced graphene into film again and thus improved the electrochemical performances. Specifically, both the capacitances in aqueous electrolytes and the rate performance in ionic liquid electrolytes were improved significantly and thereby, both high-energy and high-power density were obtained.

Graphene free-standing films are usually used directly as electrodes for various energy storage devices. In a typical synthesis process, GO sheets are vacuum filtrated into free-standing films and then chemically, thermally or laser reduced into graphene films. Fig. 1 shows a typical synthesis process of graphene films using laser reduction method. The GO sheets will stack together easily during the vacuum filtration process. Laser irradiation can simultaneously convert GO into graphene and exfoliate graphene sheets to a certain extent. The loose laser reduced graphene oxide (LrGO) can be directly pressed or re-dispersed and re-vacuum filtrated into films. However, the LrGO film will restack again and the performance will decrease accordingly. To further unstack the

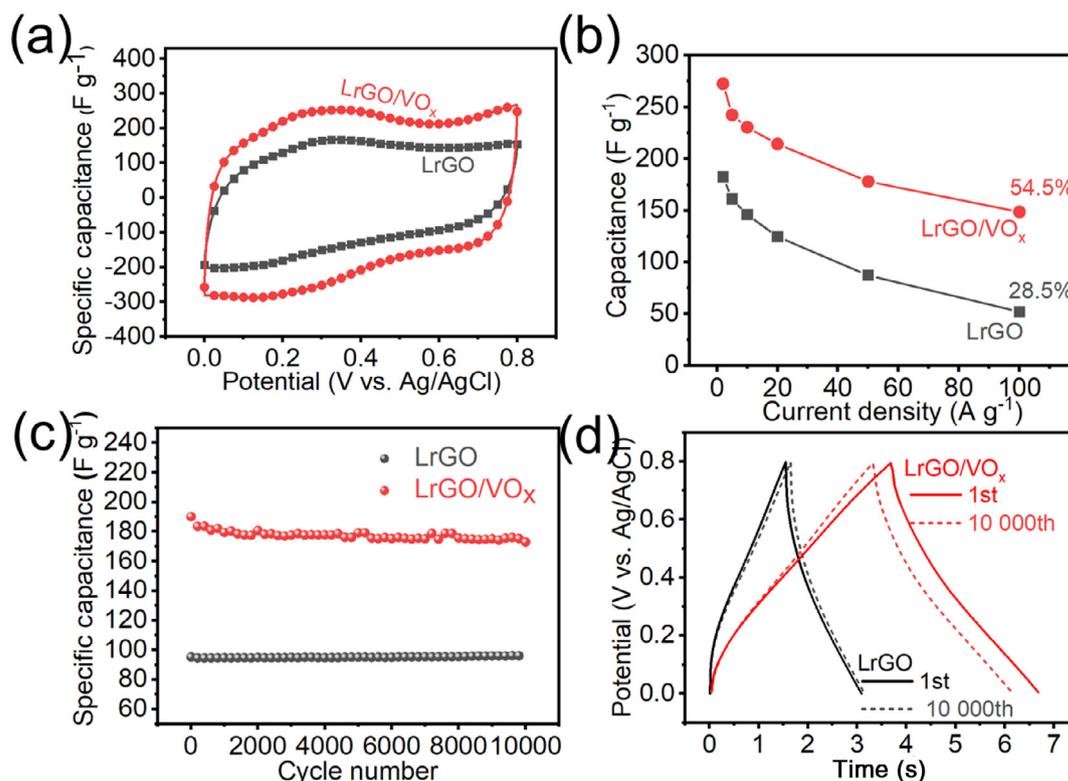
graphene sheets,  $\text{pri-VO}_x$  nanofibers are pre-inserted into GO interlayer space as a pretreatment method to prevent GO sheets from restacking. The  $\text{pri-VO}_x$  nanofibers can be converted into nanoparticles (containing  $\text{V}_2\text{O}_5$  and  $\text{VO}_2$ , Fig. S2) that chemically bonded to the graphene sheet during the laser irradiation process. When directly pressed or re-dispersed and re-vacuum filtrated into films, the  $\text{VO}_x$  nanoparticles can keep the interlayer space between adjacent graphene layers. By using this combined pretreatment and posttreatment method, highly unstacked graphene film was obtained.

It can be seen that the GO sheets are well separated by  $\text{pri-VO}_x$  nanofibers by comparing Fig. 2(a and c). The  $\text{GO/pri-VO}_x$  film shows clear layered edges, indicating the loose stacked nature of the film (compare the inset images of Fig. 2a and c). The  $\text{GO/pri-VO}_x$  film also shows increased interlayer space (Fig. 2e) because of the intercalation of  $\text{pri-VO}_x$  nanofibers. After laser reduction, the GO film pre-intercalated with  $\text{pri-VO}_x$  nanofibers ( $\text{LrGO/VO}_x$ ) becomes more loose stacked (Fig. 2d) than LrGO (Fig. 2a) and shows even higher specific surface area than LrGO (Fig. S3). The  $\text{pri-VO}_x$  nanofibers are laser converted into  $\sim 5$  nm particles on graphene sheet (inset of Fig. 2d), which are both horizontally and vertically distributed in the  $\text{LrGO/VO}_x$  film uniformly (Fig. S4). During the laser irradiation process, the GO sheets are reduced into graphene with the removal of large amount of C–O bonds (Fig. 2f and g). Surprisingly, a new peak appears at a binding energy of 530.4 eV, which can be assigned to the V–O–C bond, indicating that  $\text{VO}_x$  nanoparticles are chemically bonded to the graphene sheets.

When pressed into film, the higher SSA and larger interlayer space of  $\text{LrGO/VO}_x$  can be retained and accordingly, will contributed to higher capacitance (Fig. 3a). The capacitance of  $\text{LrGO/VO}_x$  was compared to the reported capacitances of graphene reduced by various techniques (Table S1). The  $\text{LrGO/VO}_x$  showed a high capacitance of 272.6 F/g at a current density of 2 A/g with relatively high areal mass loading of 1.2 mg/cm<sup>2</sup>. This capacitance seems to be among the highest values of reported reduced graphene oxide supercapacitors. Moreover,  $\text{LrGO/VO}_x$  film also shows higher rate performance (Fig. 3b). When charged/discharged at a very high current density of 100 A/g, 54.5% of the initial capacitance at 2 A/g can be retained for  $\text{LrGO/VO}_x$  film. While only 28.5% of the initial capacitance can be retained for the LrGO film. The  $\text{LrGO/VO}_x$  film also shows excellent cycling performance and



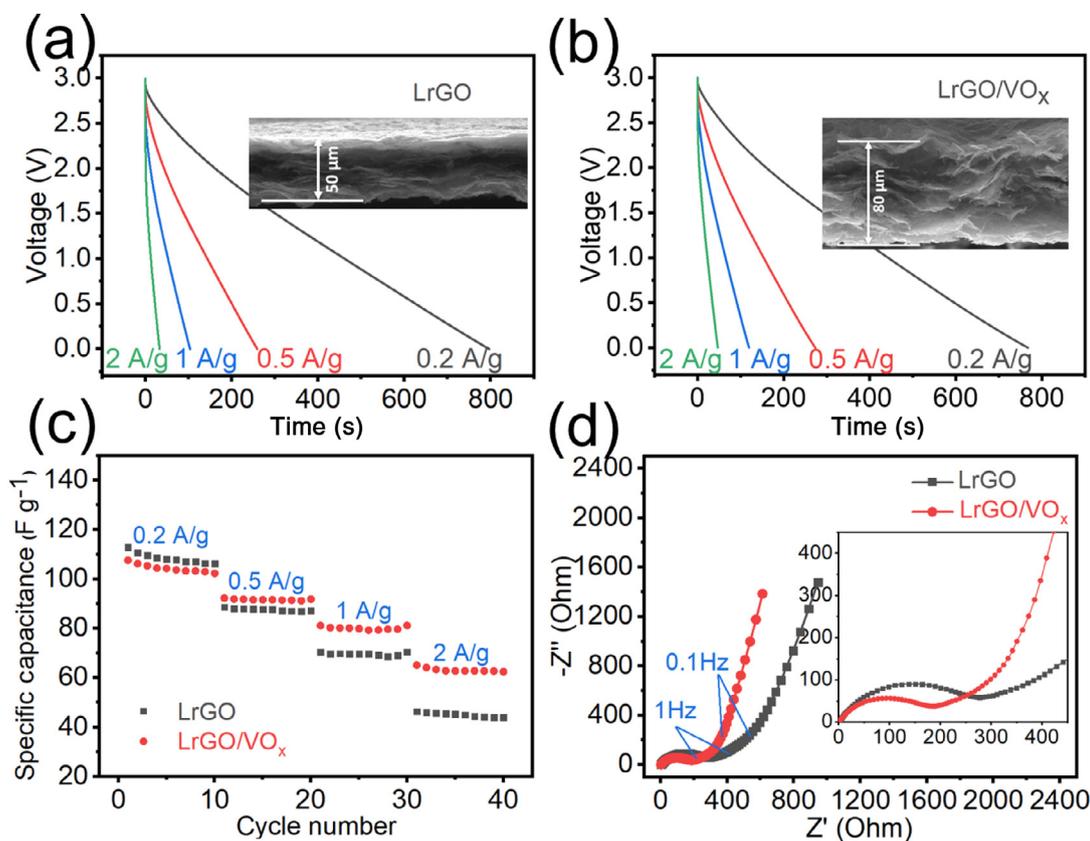
**Fig. 2.** Morphology and structure changes after laser irradiation. Scanning electron microscopy (SEM) images of GO (a) before and (b) after laser reduction; SEM images of GO/pri-VO<sub>x</sub> (c) before and (d) after laser irradiation; (e) X-ray diffraction (XRD) and (f,g) X-ray photoelectron spectroscopy (XPS) of GO and GO/pri-VO<sub>x</sub> before and after laser treatment.



**Fig. 3.** Electrochemical performance in 1 M H<sub>3</sub>PO<sub>4</sub> aqueous electrolyte. (a) Cyclic voltammetry (CV) curves of LrGO and LrGO/VO<sub>x</sub> at a scan rate of 100 mV/s. (b) Rate performance comparison of LrGO and LrGO/VO<sub>x</sub> at different current densities. (c) Cycle performance of LrGO and LrGO/VO<sub>x</sub>. (d) Charge/discharge curves of LrGO and LrGO/VO<sub>x</sub> before and after cycle performance test.

>90% of the capacitance can be retained after 10,000 cycles (Fig. 3c). By comparing the charge/discharge curve before and after cycle (Fig. 3d), it can be seen that the capacitance loss comes mainly from the high potential range for LrGO/VO<sub>x</sub>, which may be ascribed to the capacitance loss of the VO<sub>x</sub> nanoparticles. According to Fig. 3(d), the Coulombic efficiencies for LrGO before and after cycle are 90.9% and 99.4%, respectively and for LrGO/VO<sub>x</sub>, the efficiencies

are 82.4% and 85.4%, respectively. The low Coulombic efficiency of LrGO/VO<sub>x</sub> can be ascribed to the water split catalyzed by VO<sub>x</sub> since a significant increase in current density was observed around 0.8 V from the CV curve of the bare VO<sub>x</sub> (Fig. S7b, Supporting Information). To better understand the electrochemical behavior, the electrochemical performance of bare-VO<sub>x</sub> was tested as shown in Fig. S7(a). The bare VO<sub>x</sub> shows a capacitance as low as 39 F/g, which is



**Fig. 4.** Electrochemical performance in ion liquid electrolyte. Charge/discharge curves of (a) LrGO and (b) LrGO/VO<sub>x</sub> coin cells in EMI-BF<sub>4</sub> ionic liquid electrolyte within a voltage window of 3 V. (c) Rate performance of LrGO and LrGO/VO<sub>x</sub> within a voltage window of 3 V. (d) EIS at 0 V of LrGO and LrGO/VO<sub>x</sub> coin cells after rate performance test.

much lower than that of LrGO. After cycled at 50 A/g in 1 M H<sub>3</sub>PO<sub>4</sub> electrolyte for 1.5 h (corresponding to the time used for the first 1000 cycles of LrGO/VO<sub>x</sub>, the capacitance of bare VO<sub>x</sub> decreased to 21 F/g (~54% remained, nearly half lost and half remained). Therefore, assuming the capacitance fading of LrGO/VO<sub>x</sub> at the first 1000 cycles (~9%) is ascribed to the capacitance loss of VO<sub>x</sub> and thus the capacitance contribution of VO<sub>x</sub> in LrGO/VO<sub>x</sub> after 1000 cycles can be roughly calculated to be ~10%. However, the capacitance of LrGO/VO<sub>x</sub> is ~80% higher than LrGO at after 1000 cycles, indicating that the capacitance contribution from VO<sub>x</sub> is limited and can be negligible.

Symmetric LrGO and LrGO/VO<sub>x</sub> coin cells were assembled to test the electrochemical performance in EMI-BF<sub>4</sub> ionic liquid electrolyte. The discharge curves of LrGO and LrGO/VO<sub>x</sub> are shown in Fig. 4(a and b), respectively. The LrGO shows a slightly higher capacitance than that of LrGO/VO<sub>x</sub> at a low current density of 0.2 A/g. The capacitance difference is only 3.5%, which may be resulted from the inactive mass content of VO<sub>x</sub> in LrGO/VO<sub>x</sub> (~3.7 wt%, Fig. S1). The LrGO/VO<sub>x</sub> with enlarged interlayer space (insets of Fig. 4a and b) shows higher rate performance (Fig. 4c). The capacitance retention increases from 46.3% to 65.1% when the charge/discharge current density increases from 0.2 to 2 A/g within a voltage window of 3 V. The higher rate performance of LrGO/VO<sub>x</sub> calculated from the discharge curves showed consistency with the CV results shown in Fig. S8, where larger CV area was obtained for LrGO/VO<sub>x</sub>. The 4 V performance has also been tested and is shown in Fig. S9, which also displays a rate capability increase. The high rate performance may be ascribed to the increased surface area and the decreased resistance at lower frequencies (Figs. 4d and S9d). The increased rate capability may be also attributed to the loose stacked structure of LrGO/VO<sub>x</sub> (the inset SEM images

in Fig. 4a and b) where VO<sub>x</sub> works as spacer between LrGO sheets and thus electrolyte ions can transport faster. As shown in Fig. S11, the 4 V voltage window allows the LrGO/VO<sub>x</sub> to achieve a high energy density of 84.4 Wh/kg at 0.5 A/g. Even at a high power density of 1000 W/kg, the LrGO/VO<sub>x</sub> shows a high energy density of 47.3 Wh/kg. Fig. 4(d) shows that the impedance values at 10<sup>5</sup> Hz are almost the same for both supercapacitors, indicating the inner resistance of graphene has not been significantly changed by VO<sub>x</sub>. However, the semi-circle diameter of the EIS spectra of LrGO/VO<sub>x</sub> supercapacitor at high frequencies are smaller and the slope at lower frequencies are larger than LrGO supercapacitor, which implies that both the charge transfer and diffusion resistances are smaller for LrGO/VO<sub>x</sub> supercapacitor compared to LrGO supercapacitor.

In conclusion, pre-inserted pri-VO<sub>x</sub> nanofibers can effectively prevent the GO sheets from restacking together. Laser reduction is an effective method for restacked GO film reduction and exfoliating. By combining these two methods, the GO sheets pre-exfoliated by pri-VO<sub>x</sub> inserting can be further exfoliated by laser irradiation and highly unstacked graphene can be obtained. Consequently, ultrahigh capacitance of 272.6 F/g of the graphene is obtained in 1 M H<sub>3</sub>PO<sub>4</sub> electrolyte. This value is among the highest values that have been reported in graphene-based supercapacitors using aqueous electrolytes. An improved rate performance has also been obtained in neat ionic liquid electrolyte for the VO<sub>x</sub>-bonded graphene.

#### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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## Supplementary materials

Supplementary material associated with this article can be found, in the online version, at doi:10.1016/j.jechem.2020.02.015.

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